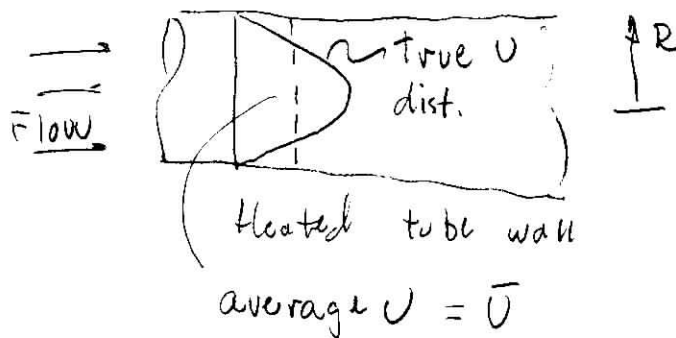
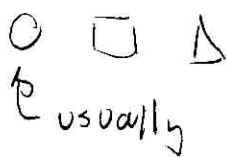


Internal Flows

stationary tube wall

10

Ducts



In fully developed region (far down pipe)

Recall $\dot{m} = \int_{A_c} \rho U(r) dA_c = \rho \bar{U} A_c$

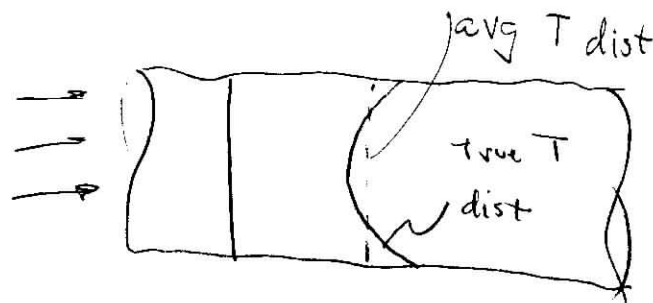
or
$$\bar{U} = \frac{\int_{A_c} \rho U(r) dA_c}{\rho A_c} = \frac{\int_{r=0}^R \rho U(r) 2\pi r dr}{\rho \pi R^2} = \frac{2}{R^2} \int_{r=0}^R U(r) r dr$$

gotta know

For energy transport down pipe

$$\dot{E}_{fluid} = \dot{m} c_p \bar{T}$$

$$= \int_{A_c} \rho c_p T(r) U(r) dA_c$$



so
$$\bar{T} = \frac{\int_{r=0}^R \rho c_p T(r) U(r) 2\pi r dr}{\rho \bar{U} (\pi R^2) c_p} = \frac{\int \dot{m} c_p}{\int \dot{m} c_p} = \frac{2}{\bar{U} R^2} \int_{r=0}^R T(r) U(r) r dr$$

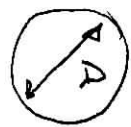
usually evaluate mtl. properties @ a bulk temperature (\bar{T})

$$T_{bulk} = \frac{T_{inlet} + T_{exit}}{2} = \frac{T_i + T_e}{2}$$

hydraulic Diameter

In a circular tube

$$Re = \frac{\bar{U} D}{\nu} = \frac{\rho \bar{U} D}{\mu} = \frac{4 \dot{m}}{\mu \pi D}$$



Recall

$$\nu = \frac{\mu}{\rho}$$

We call the Hydraulic dia

$$D_{hyd} \equiv \frac{4 A_{cross}}{perim.}$$

so for a cyl.

$$D_h = \frac{4 \pi \left(\frac{D^2}{4}\right)}{\pi D} = D$$

Usually for a circ. tube

$$Re < 2000$$

$$Re \sim 4000$$

$$Re > 10,000$$

prob. lam.



Transition

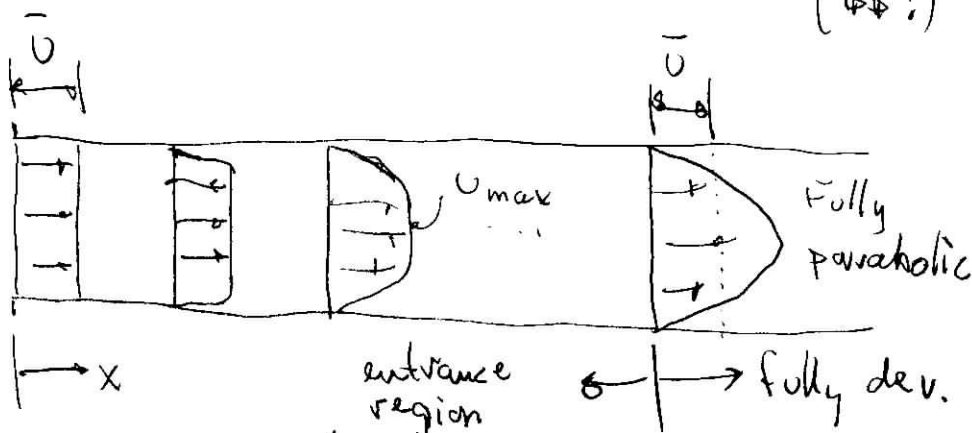
lam. to Turb.



fully turb.

Special cases: ... lam. flow has been maintained upto $Re \sim 100,000$ (\$\$!!)

Entrance Region



Note \bar{U} is \downarrow as $x \uparrow$

but $U_{max} \uparrow$ as $x \uparrow$

Also eventually δ_m grows from 0 at entrance to $\delta_m = R$ at fully dev. conditions (not considered entrance region any more).

Similar results for thermal $\delta_T(x)$

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once we hit the end of the entrance region (we are in the fully developed region) we have

$$\frac{\partial u(r, x)}{\partial x} = 0 \Rightarrow u = u(r) \text{ only}$$

$$\frac{\partial}{\partial x} \left[\frac{T_s(x) - T(r, x)}{T_s - \bar{T}} \right] = 0$$

$$\Rightarrow T = T(r) \text{ only}$$

Conseq. of
 $T = T(r) \text{ only}$

$$\text{AND } \left. \frac{\partial T}{\partial r} \right|_{r=R} \neq f(x)$$

$$\text{so } h(x) \equiv \frac{k \left(\left. \frac{\partial T}{\partial r} \right|_{r=R} \right)}{T_s - \bar{T}} \neq f(x)$$

(local value)

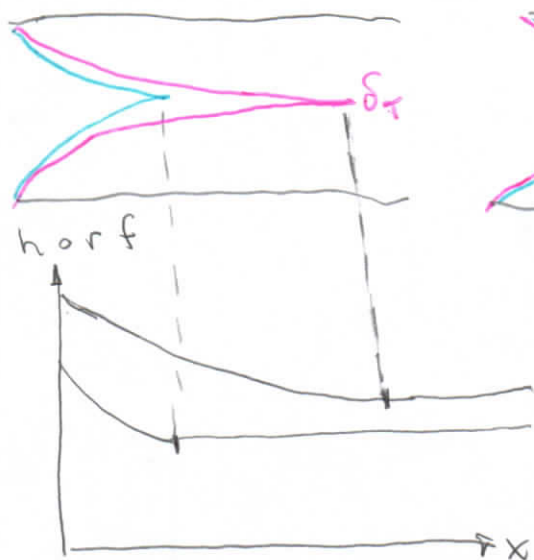
This in the fully developed section

fric factor $f \neq \text{function}(x)$

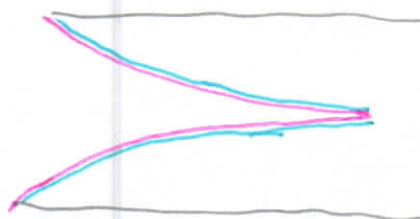
$h \neq \text{function}(x)$

Role of $Pr = \nu/\alpha$

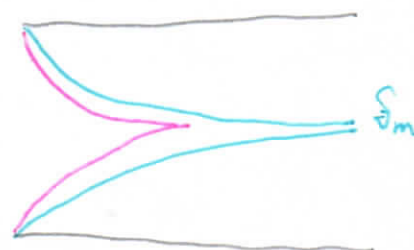
$Pr \gg 1$



$Pr = 1$



$Pr < 1$



So what are entry length results?

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Laminar
flow

$$L_{\text{hyd.}} \approx 0.05 \text{ Re } D$$

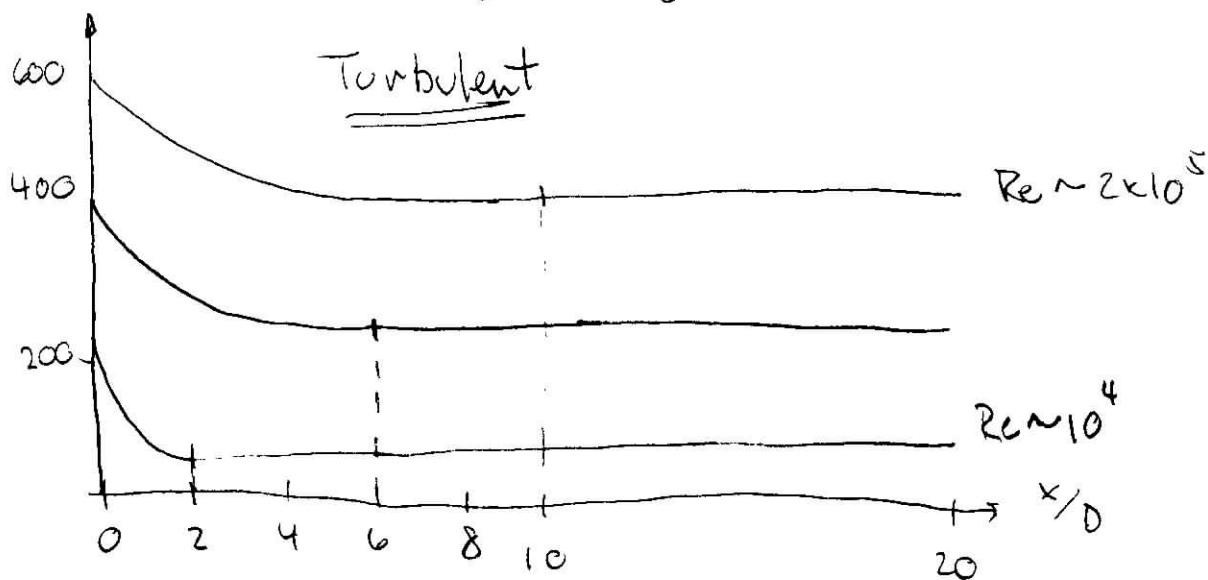
$$L_{\text{therm.}} \approx 0.05 \text{ Re } D \text{ Pr}$$

- Thus if $\text{Re} = 20$ the entry length \approx diameter of tube
if $\text{Re} \approx 2300$ $\approx 111 D$

Turb,
flow

$$L_h \approx L_t \approx 10 D$$

$\text{Nu}(x)$ for either $T_s = \text{const}$ or $q_s = \text{const}$ at surface

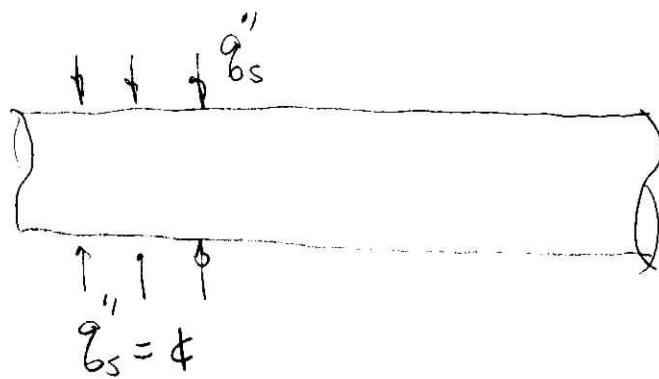


Observations : • $\text{Nu} \uparrow$ in entrance region

- Nu are all $\approx \text{const}$ for $\frac{x}{D} > 10$ (no matter the Re)
 - Nu is almost indep. of B.C. ($T_s = \text{const}$ or $q_s = \text{const}$) (int. or fully dev.)
- \Rightarrow Might as well just use turb. correlations for either type of B.C.

Constant heat
Flux at
Surface

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Surf.
heat
flux

$$q''_s = h(x) (T_s(x) - \bar{T}(x)) \quad \left[\frac{\text{W}}{\text{m}^2} \right]$$

\uparrow local heat transf. coef. \uparrow \bar{T} changes down the pipe

There can still be lots of things that vary with x on RHS.

$$\dot{Q} = q''_s A_s = \dot{m} C_p (T_{\text{exit}} - T_{\text{inlet}}) \quad [\text{W}]$$

$$\text{Then } T_{\text{exit}} = T_{\text{in}} + \frac{q''_s A_s}{\dot{m} C_p}$$

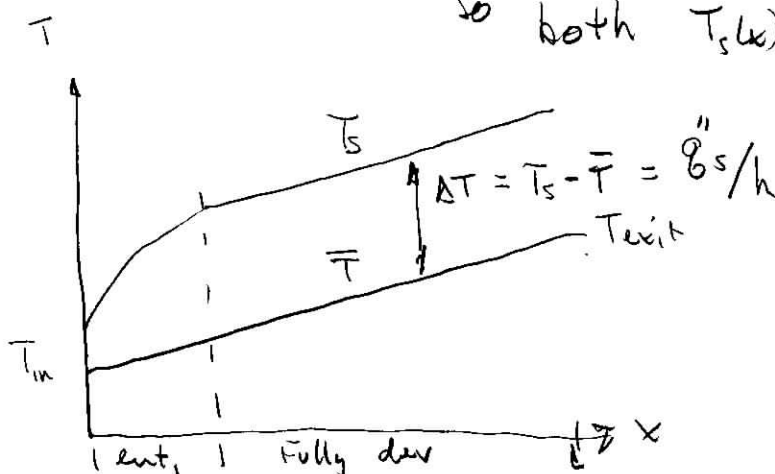
$\Rightarrow \bar{T}$ increases linearly in dir. of flow

recall $A_s = \text{surf. area}$
 $= \text{perim} \cdot \text{length}$

But also $q''_s = h(T_s(x) - \bar{T}(x)) \Rightarrow T_s = \bar{T} + \frac{q''_{\text{surf}}}{h}$

But in fully dev. region $h = \text{const}$ (and $q''_s = \text{const}$)

so both $T_s(x)$ and $\bar{T}(x)$ increase linearly



In fully developed region, the dimensionless $\Theta = \frac{T_s - T}{T_s - \bar{T}}$ is \neq wrt x 8/10

$$\text{so } \frac{\partial}{\partial x} \left(\frac{T_s - T}{T_s - \bar{T}} \right) = 0 = \frac{\frac{\partial T_s}{\partial x} - \frac{\partial T}{\partial x}}{(T_s - \bar{T})} \quad \text{leads to } \frac{\partial T}{\partial x} = \frac{\partial T_s}{\partial x} = \frac{dT_s}{dx}$$

in which case

$$\frac{\partial T}{\partial x} = \frac{dT_s}{dx} = \frac{dT_m}{dx} = \frac{q_s'' \cdot \text{Perim}}{\dot{m} C_p} = \neq \quad (*)$$

from simple
heat bal. of
a slice through
the pipe

None of these gradients (wrt x) vary
down the length of the pipe.

② (*) from $x=0$ where $\bar{T} = T_{\text{inlet}}$ we get

$$\bar{T} = T_{\text{inlet}} + \left(\frac{q_s'' P}{\dot{m} C_p} \right) x$$

Eval. this @ $x=L$ (where $\bar{T} = T_{\text{exit}}$) and recall $A_s = P \cdot L$

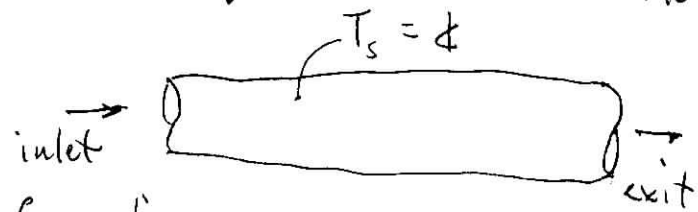
$$\text{we get } T_{\text{exit}} = T_{\text{inlet}} + \frac{q_s'' A_s}{\dot{m} C_p}$$

For a circular tube $P = 2\pi R$ $\dot{m} = \rho \bar{U} A_{\text{cross}} = \rho \bar{U} (\pi R^2)$

$$\frac{\partial T}{\partial x} = \frac{dT_s}{dx} = \frac{d\bar{T}}{dx} = \frac{2q_s''}{\rho \bar{U} C_p R} = \neq \quad \text{Const. } q_s'' \text{ flux @ tube wall.}$$

Now what about B.C. of $T_s = \text{const}$ @ wall?

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Roughly speaking, from Newton's law of cooling

$$\dot{q} = h A_{\text{surf}} \Delta T_{\text{avg}} = h A_s (\overline{T}_s - \overline{T}) \quad [W]$$

$\left\{ \begin{array}{l} \text{avg} \\ \text{value over} \\ \text{surface} \end{array} \right.$
 $\left\{ \begin{array}{l} \text{average} \\ \text{spatial} \\ \dots \text{ whatever} \\ \text{that is} \end{array} \right.$

There are 2 ways of thinking about this,

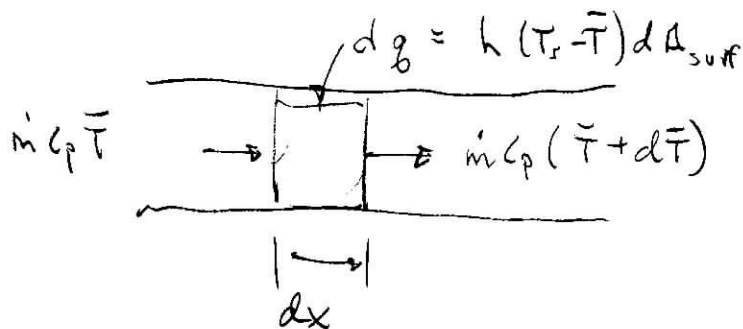
Arithmetic mean temp. difference $\Delta T_{\text{a.m.}}$

$$\begin{aligned} \Delta T_{\text{avg}} &\doteq \Delta T_{\text{a.m.}} = \frac{\Delta T_i + \Delta T_e}{2} = \frac{(\overline{T}_s - T_i) + (\overline{T}_s - T_e)}{2} \\ &= \overline{T}_s - \left(\frac{T_i + T_e}{2} \right) \\ &= \overline{T}_s - \overline{T}_{\text{bulk}} \end{aligned}$$

arith. average
of temp dif.
btwn surf and
fluid at inlet
and exit

leads to reasonable results, but often fails.

Consider an energy balance on



$$\dot{m} c_p d\overline{T} = h (\overline{T}_s - \overline{T}) dA_{\text{surf}}$$

note $dA_{\text{surf}} = p \cdot dx$
 \uparrow perim.

and $d\overline{T} = -d(\overline{T}_s - \overline{T})$

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leads to $\frac{d(T_s - \bar{T})}{(T_s - \bar{T})} = -\frac{h p}{\dot{m} c_p} dx$ \int from $x=0$ ($\bar{T} = T_i$) to $x=L$ ($\bar{T} = T_e$)

leads to $\ln \left(\frac{T_s - T_e}{T_s - T_i} \right) = -\frac{h A_s}{\dot{m} c_p}$ ** (h is the ϕ average heat trans. coeff.)

exponentiate to get

$$\frac{T_e - T_s}{T_i - T_s} = \exp \left(\frac{-h A_s}{\dot{m} c_p} \right) \quad \text{where } A_s = pL$$

If you replace $A_s = pL$ by pX , you get soln. for

$$\frac{\bar{T}(x) - T_s}{T_i - T_s} = \exp \left(\frac{-h p x}{\dot{m} c_p} \right)$$

We define $\frac{h A_s}{\dot{m} c_p} = NTU$ Number of transfer units

If $NTU > 5$ then $T_{exit} \approx T_{surf}$

$NTU > 5$ no more heat trans. happens. No need to add more length.

$NTU < 5$ helps to add more length to pipe

Solve q for $\dot{m} c_p = \frac{-h A_s}{\ln \left[\frac{T_s - T_e}{T_s - T_i} \right]}$

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put this into $q = \dot{m} c_p (T_e - T_i) \text{ [W]}$ near start of today

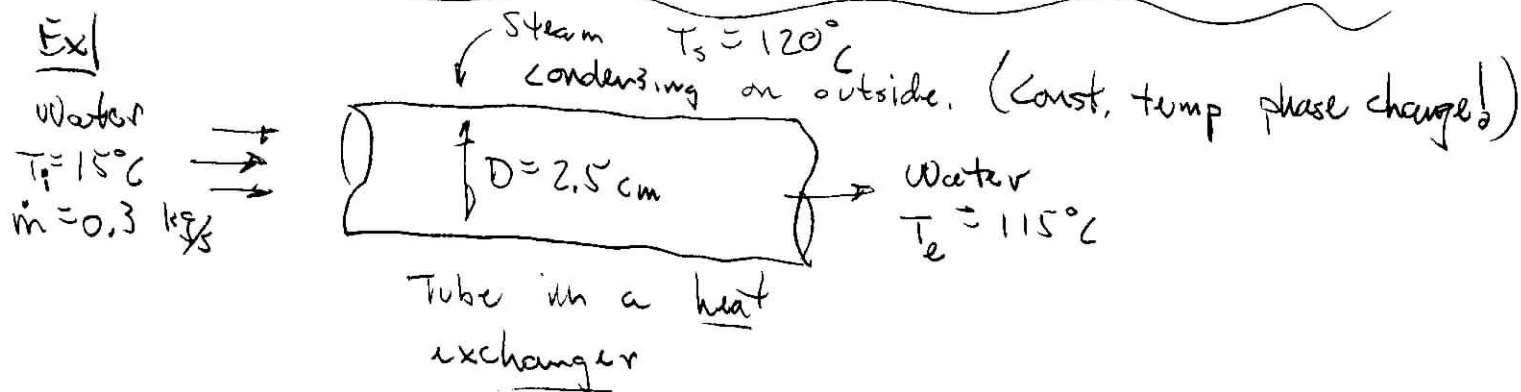
we get

$$q = h A_s \left\{ \frac{(T_i - T_e)}{\ln \left[\frac{T_s - T_e}{T_s - T_i} \right]} \right\} = h A_s \Delta T_{lm}$$

where $\Delta T_{\log \text{ mean}} = \frac{T_i - T_e}{\ln \left(\frac{T_s - T_e}{T_s - T_i} \right)} = \frac{\Delta T_e - \Delta T_i}{\ln \left(\frac{\Delta T_e}{\Delta T_i} \right)}$

Log mean
Temp. Dif.

This is the way to go rather than the Arit. mean temp. dif. method.



If Average $h = 800 \text{ W/m}^2\text{K}$, determine tube length L for exit temp. of 115°C .

- Assumptions:
- Steady state
 - h is ∞
 - R_{copper} is negl.
(inside tube $T = \text{outside tube } T$)
 $k_{\text{cu}} \uparrow \uparrow$

Math. prop: $\odot T_{\text{bulk}} = \frac{(15 + 115)^{\circ}\text{C}}{2} = 65^{\circ}\text{C}$

$c_p = 4187 \frac{\text{J}}{\text{kg} \cdot \text{K}}$

phase change
heat @ 120°C
release $= 2203 \frac{\text{kJ}}{\text{kg}}$

Let's go:

for the water in the tube

$\dot{q} = \dot{m} c_p (T_e - T_i) = \dots = 125.6 \text{ kW}$
+ trans.

To calc ΔT_{lm} we need $\Delta T_e = T_s - T_e = 120^{\circ}\text{C} - 115^{\circ}\text{C} = 5^{\circ}\text{C}$

$\Delta T_i = T_s - T_i = 120^{\circ}\text{C} - 15^{\circ}\text{C} = 105^{\circ}\text{C}$

$\Delta T_{\text{lm}} = \frac{\Delta T_e - \Delta T_i}{\ln\left(\frac{\Delta T_e}{\Delta T_i}\right)} = \dots = 32.85^{\circ}\text{C}$

The heat transfer surface area is from

$\dot{q} = h A_s \Delta T_{\text{lm}}$ or $A_s = \frac{\dot{q}}{h \Delta T_{\text{lm}}} = \dots = 4.78 \text{ m}^2$

But

$A_s = \pi D L$ so $L = \frac{A_s}{\pi D} = 61 \text{ m}$

Comments $T_{\text{bulk}} = 65^{\circ}\text{C}$ so arithmetic $\Delta T_{\text{am}} = (120 - 65)^{\circ}\text{C} = 55^{\circ}\text{C}$
which would give $L = 36 \text{ m}$ which is really wrong

Use ΔT_{lm}

